# $B \rightarrow D^{*}$ Form Factors from Light-Cone Sum Rules 

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## Motivations: why do we need B to D* FFs?

- $\left|V_{c b}\right|$ extraction from branching ratios of $B \rightarrow D^{*} \mu \nu$
- prediction of $R_{D^{*}}$ in the SM , i.e. to constrain NP contributions to $b \rightarrow c l \bar{\nu}$
- LCSRs complement Lattice results and Heavy Quark Expansion relations used in present analyses
- B-LCSRs have $1 / m_{b}$ corrections (related to twist expansion), but there is no $1 / m_{c}$ expansion!
- we present new twist 4 corrections to the $B \rightarrow D^{*}$ LCSRs, higher twists are expected to give corrections only of the order $O\left(1 / m_{b}^{2}\right)$
- $O\left(\alpha_{s}\right)$ corrections are not considered


## Light-Cone Sum Rules in a nutshell

- determine products of exclusive hadronic matrix elements from an artificial, less-exclusive, non-local hadronic matrix element $\Pi\left(k^{2}, q^{2}\right)$
- $\Pi\left(k^{2}, q^{2}\right)$ calculable for kinematics that impose light-cone dominance of the non-local operator
- results

$$
\Pi\left(k^{2}, q^{2}\right)=f_{B} m_{B} \int d s \sum_{n, t} \frac{J_{n, t}\left(s, q^{2}\right)}{\left[k^{2}-s\right]^{n}} \phi_{t}(s)
$$

- $J_{n, t}$ can be computed from a hard scattering kernel
- B-meson Light-Cone Distribution Amplitudes (LCDAs) $\phi_{t}$ are necessary non-perturbative input
- general $B \rightarrow V, B \rightarrow P$ results available
- new insights on LCDAs triggered our revisiting of these sum rule results


## Preliminary Results and Comparison

|  | FKKM2008 | GKvD2018 |  |  |
| :---: | :---: | :---: | :---: | :---: |
|  |  | NEW Contrib. |  |  |
| B $\rightarrow \mathbf{D}^{*}$ FF | 2pt tw2+3 +3pt | 2pt tw2+3 | 2pt tw4 | 3pt tw3+4 |
| $A_{1}\left(q^{2}=0\right)$ | 0.73 | 0.65 | -0.11 | $?$ |
| $A_{2}\left(q^{2}=0\right)$ | 0.66 | 0.57 | -0.21 | $?$ |
| $A_{0}\left(q^{2}=0\right)$ | 0.78 | 0.70 | -0.01 | $?$ |
| $A_{0}(0) / A_{1}(0)$ | 1.07 | 1.08 | +0.21 | $?$ |

[using the same input parameters, with $q^{2}$ the dilepton mass square]
$\phi_{+}, \phi_{-}$2-particle L+NL twist contributions [Faller/Khodjamirian/Klein/Mannel '06] $\mathbf{g}_{+}$new 2-particle NNL twist contributions [Gubernari/Kokulu/van Dykw.i.p.] $\phi_{3}, \phi_{4}$ new and self-consistent 3-particle NL+NNL twist contr.
[Gubernari/Kokulu/van Dyk w.i.p.]

## Plans for presentation of results

- we plan to give numerical results for all form factors at $q^{2}=0$ and $q^{2}=-5 \mathrm{GeV}^{2}$
- we consider $q^{2}=+5 \mathrm{GeV}^{2}$ as an additional point, but we will check convergence of the twist expansion first before committing to use it
- we plan to provide correlation matrices across form factors and across $q^{2}$
- we plan to provide numerical results in machine-readable form - probably JSON/YAML files, similar to what has been done for light-meson LCSRs
- numerical evaluations are carried out with EOS and the code will be made publicly available at https://github.com/eos/eos


## Backup slides

## Power corrections

- correlator is calculated with on-shell $B$ meson, using its Light-Cone Distribution Amplitudes (LCDAs)
- $B$-meson LCDAs are defined for bi-local currents involving an HQET field $h_{v}$
- power corrections to this involve power of the covariant derivative $i D^{\mu}$
- strings of the type $i D^{\mu_{1}} i D^{\mu_{2}} \ldots i D^{\mu_{n}}$ are incorporated in LCDAs of increasing (collinear) twist


## Benefits of the Braun et al. basis

- $\phi_{3}, \phi_{4}, \ldots$ are LCDAs of definite collinear twist $3,4, \ldots$
- LCDAs of twists $\geq 5$ are expected to contribute beyond the next-to-leading $1 / m_{b}$ corrections!
- inserting a gluon field adds at least one unit of twist
- 2-particle LCDAs start at twist 2, and are included in our results (up to and including twist 4)
- 3-particle LCDAs start at twist 3, and are included in our results (up to and including twist 4)
- 4-particle LCDAs start at twist 4, and are not included in our results
- 4-particle LCDAs have autonomous RG behaviour, do not mix with 3-particle LCDAs

