# $B \rightarrow D^*$ Form Factors from Light-Cone Sum Rules



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Challenges in Semileptonic *B* Decays MITP, Mainz



## Motivations: why do we need B to D\* FFs?

- $|V_{cb}|$  extraction from branching ratios of  $B 
  ightarrow D^* \mu 
  u$
- prediction of  $R_{D^*}$  in the SM, i.e. to constrain NP contributions to  $b \rightarrow c l \bar{\nu}$
- LCSRs complement Lattice results and Heavy Quark Expansion relations used in present analyses
- B-LCSRs have 1/m<sub>b</sub> corrections (related to twist expansion), but there is no 1/m<sub>c</sub> expansion!
- we present new twist 4 corrections to the B → D\* LCSRs, higher twists are expected to give corrections only of the order O(1/m<sup>2</sup><sub>b</sub>)
- O(α<sub>s</sub>) corrections are not considered

# Light-Cone Sum Rules in a nutshell

- determine products of exclusive hadronic matrix elements from an artificial, less-exclusive, non-local hadronic matrix element Π(k<sup>2</sup>, q<sup>2</sup>)
- Π(k<sup>2</sup>, q<sup>2</sup>) calculable for kinematics that impose light-cone dominance of the non-local operator
- results

$$\Pi(k^{2}, q^{2}) = f_{B}m_{B} \int ds \sum_{n,t} \frac{J_{n,t}(s, q^{2})}{[k^{2} - s]^{n}} \phi_{t}(s)$$

- $J_{n,t}$  can be computed from a hard scattering kernel
- B-meson Light-Cone Distribution Amplitudes (LCDAs)  $\phi_t$  are necessary non-perturbative input
  - general  $B \rightarrow V$ ,  $B \rightarrow P$  results available [Khodjamirian et al. '06 + '08]
  - new insights on LCDAs triggered our revisiting of these sum rule results [Braun/Ji/Manashov '17]

	FKKM2008	GKvD2018		
			NEW Contrib.	
$\textbf{B} \rightarrow \textbf{D}^* \; \textbf{F}\textbf{F}$	2pt tw2+3 +3pt	2pt tw2+3	2pt tw4	3pt tw3+4
$A_1(q^2 = 0)$	0.73	0.65	-0.11	?
$A_2(q^2 = 0)$	0.66	0.57	-0.21	?
$A_0(q^2 = 0)$	0.78	0.70	-0.01	?
$A_0(0)/A_1(0)$	1.07	1.08	+0.21	?

[using the same input parameters, with  $q^2$  the dilepton mass square]

 $\phi_{+}, \phi_{-}$  2-particle L+NL twist contributions [Faller/Khodjamirian/Klein/Mannel '06]  $\mathbf{g}_{+}$  new 2-particle NNL twist contributions [Gubernari/Kokulu/van Dyk w.i.p.]  $\phi_{3}, \phi_{4}$  new and self-consistent 3-particle NL+NNL twist contr. [Gubernari/Kokulu/van Dyk w.i.p.] 3/4

- we plan to give numerical results for all form factors at  $q^2 = 0$  and  $q^2 = -5 \text{ GeV}^2$
- we consider  $q^2 = +5 \text{ GeV}^2$  as an additional point, but we will check convergence of the twist expansion first before committing to use it
- we plan to provide correlation matrices across form factors and across q<sup>2</sup>
- we plan to provide numerical results in machine-readable form
  - probably JSON/YAML files, similar to what has been done for light-meson LCSRs [Bharucha/Straub/Zwicky '15]
- numerical evaluations are carried out with EOS and the code will be made publicly available at https://github.com/eos/eos

# Backup slides

- correlator is calculated with on-shell *B* meson, using its Light-Cone Distribution Amplitudes (LCDAs)
- *B*-meson LCDAs are defined for bi-local currents involving an HQET field  $h_v$
- power corrections to this involve power of the covariant derivative  $iD^{\mu}$
- strings of the type *iD<sup>µ1</sup> iD<sup>µ2</sup> ... iD<sup>µn</sup>* are incorporated in LCDAs of increasing (collinear) twist

- $\phi_3$ ,  $\phi_4$ , ... are LCDAs of definite collinear twist 3, 4, ...
- LCDAs of twists  $\geq$  5 are expected to contribute beyond the next-to-leading  $1/m_b$  corrections! [Braun/Ji/Manashov '17]
- inserting a gluon field adds at least one unit of twist
  - 2-particle LCDAs start at twist 2, and are included in our results (up to and including twist 4)
  - 3-particle LCDAs start at twist 3, and are included in our results (up to and including twist 4)
  - 4-particle LCDAs start at twist 4, and are not included in our results
  - 4-particle LCDAs have autonomous RG behaviour, *do not mix with 3-particle LCDAs*

[Braun/Ji/Manashov '17]